Quantum Field Theory

Homework #3

Hand-in time and place (week 7):

Class	Hand-in time	Hand-in place	Teaching Assistant
Tuesday 15.30-17.00	Sunday 6pm	Mathematics [†]	Johan Henriksson
Friday 14.30-16.00	Tuesday noon	Mathematics [†]	Johan Henriksson
Thursday 8.30-10.00	Monday 6pm	Mathematics [‡]	Matteo Parisi
Friday 8.30-10.00	Monday 6pm	Mathematics [‡]	Matteo Parisi

† Mezzanine level in the Mathematical Institute (Andrew Wiles Building, Woodstock Road)

DO NOT FORGET TO PUT THE NAME OF YOUR TEACHING ASSISTANT ON THE SHEET.

Useful formulae:

• Schwinger parametrisation:

$$\frac{1}{A^x} = \frac{1}{\Gamma(x)} \int_0^\infty dt \, t^{x-1} e^{-At} \,. \tag{1}$$

• Momentum integrals in d dimensions

$$\int \frac{d^d k}{(2\pi)^d} e^{-\alpha k^2} = (4\pi\alpha)^{-d/2} \,. \tag{2}$$

• Miscellaneous useful integrals:

$$\int_{0}^{1} dx \int_{0}^{1} dy \int_{0}^{1} dz \, \delta(x+y+z-1) \frac{xyz}{(xy+xz+yz)^{3}} = \frac{1}{2}, \quad (3)$$

$$\int_{0}^{1} dx \int_{0}^{1} dy \int_{0}^{1} dz \, \delta(x+y+z-1) \frac{xyz}{(xy+xz+yz)^{3}} \times$$

$$\times \log \frac{(xy+xz+yz)^{3}}{(xy+xz+yz-xyz)^{2}} = -\frac{3}{4}. \quad (4)$$

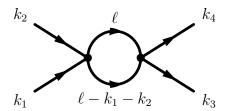
8. Feynman parametrisation

Using induction in n, prove the following formula

$$\frac{1}{a_1 \cdots a_n} = (n-1)! \int_0^1 \prod_{i=1}^n dx_i \frac{\delta(\sum_{j=1}^n x_j - 1)}{(a_1 x_1 + \dots + a_n x_n)^n}.$$
 (5)

9. Evaluating Feynman integrals

In this homework we evaluate a contribution to the four-point correlator in the ϕ^4 theory in Minkowski space at one-loop, given by the diagram



- (a) Write the dimensionless coupling constant g_0 in $4-2\epsilon$ dimensions as the bare coupling λ_0 multiplied by a suitable power of the mass scale μ .
- (b) Show that using the Feynman rules in momentum space and after Wick rotation the diagram yield

$$A(\epsilon) = \frac{g_0^2}{2} \mu^{4\epsilon} \int \frac{d^{4-2\epsilon}\ell}{(2\pi)^{4-2\epsilon}} \frac{1}{\ell^2 + m^2} \frac{1}{(\ell - k_1 - k_2)^2 + m^2}.$$
 (6)

(c) Use Feynman's identity to rewrite the loop integration as

$$A(\epsilon) = \frac{g_0^2}{2} \mu^{4\epsilon} \int_0^1 dx \int \frac{d^{4-2\epsilon}\tilde{\ell}}{(2\pi)^{4-2\epsilon}} \frac{1}{(\tilde{\ell}^2 + x(1-x)(k_1+k_2)^2 + m^2)^2}.$$
(7)

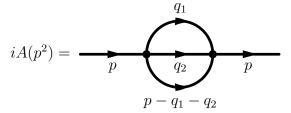
(d) Show that the momentum integral evaluates to

$$A(\epsilon) = \mu^{2\epsilon} \frac{g_0^2}{32\pi^2} \left[\frac{1}{\epsilon} - \gamma_E - \int_0^1 dx \ln\left(\frac{m^2 + (k_1 + k_2)^2 x (1 - x)}{4\pi\mu^2}\right) + \mathcal{O}(\epsilon) \right], \tag{8}$$

where γ_E is the Euler constant.

10. Wave-function renormalisation

In order to find the leading correction to wave-function renormalization constant in $\lambda_0 \phi^4$ theory in Minkowski space we will evaluate the so-called sunrise diagram:



- (a) Write down the Feynman integral for sunrise diagram in four dimensions. Remember to include the proper symmetry factor.
- (b) Use Feynman parametrisation and rewrite the Feynman integral as

$$A(p^{2}) = C \int_{0}^{1} dx \int_{0}^{1} dy \int_{0}^{1} dz \, \delta(x+y+z-1) \times \int \frac{d^{4}q_{1}}{(2\pi)^{4}} \int \frac{d^{4}q_{2}}{(2\pi)^{4}} \frac{1}{(\mathcal{D}(x,y,z,q_{1}^{2},q_{2}^{2},p^{2}))^{3}}, \quad (9)$$

where C is a constant and the function \mathcal{D} depends only on the squares of momenta q_1 and q_2 .

(c) In order to find the contribution to wave-function renormalization we need to evaluate the derivative of $A(p^2)$ with respect to p^2 . Working in Euclidean space and using dimensional regularization in $d=4-2\epsilon$ dimensions show that

$$\frac{dA(p^2)}{dp^2} = \int_0^1 dx \int_0^1 dy \int_0^1 dz \, \delta(x+y+z-1) F(x,y,z) \times \left\{ \frac{2}{\epsilon} + const + \log G(x,y,z,p^2/m^2) \right\}, \tag{10}$$

for some functions F and G

(d) Using the explicit form of this functions and working in the leading order perturbative expansion in dimensionless coupling evaluate

$$\frac{dA(p^2)}{dp^2}\Big|_{p^2=m^2} = -\frac{g_0^2}{6144\pi^4} \left(\frac{1}{\epsilon} + 2\log\frac{\mu^2}{m^2} + 2\log(4\pi) - 2\gamma_E - \frac{3}{2} + \mathcal{O}(\epsilon)\right).$$
(11)

Some of the integrals given in the introduction can be useful here.

(e) Evaluate the anomalous dimension for the field ϕ from the formula

$$\gamma(g_0) = \mu \frac{\partial}{\partial \mu} \log Z_{\phi}, \qquad (12)$$

where up the order we are working at

$$Z_{\phi}^{-1} = \frac{1}{1 + \frac{dA(p^2)}{dp^2} \Big|_{p^2 = m^2}}.$$
 (13)

Compare the result with the one found in lecture.