## General Relativity I - Problem Sheet 3

Please return solutions (collection boxes in the basement of the Mathematical Institute) no later than 15 th of November, 1 pm .

## 1. Independent Components of The Riemann Tensor

The Riemann tensor $R_{a b c}{ }^{d}$ associated to any covariant derivative $\nabla_{a}$ obeys the following algebraic identities:

1. $R_{a b c}{ }^{d}=-R_{b a c}{ }^{d}$
2. $R_{a b c}{ }^{d}+R_{b c a}{ }^{d}+R_{c a b}{ }^{d}=0$

What is the additional algebraic identity obeyed when $\nabla_{a}$ is the Levi-Civita convariant derivative assoicated to a metric $g_{a b}$ ?

In this case, show that the Riemann tensor in four dimensions has, in general, 20 independent components.

Optional: How many independent components does it have in $D$ dimensions?

## 2. Riemann tensor in two dimensions

1. Prove that in two dimensions, the Riemann tensor has the form

$$
R_{a b c d}=\frac{1}{2} R\left(g_{a c} g_{b d}-g_{a d} g_{b c}\right) .
$$

where $R$ is the Ricci scalar. Hint: express the Ricci scalar $R$ in terms of the one independent component of the Riemann tensor, say $R_{1212}$.
2. Show that Einstein tensor $G_{a b}=R_{a b}-\frac{1}{2} g_{a b} R$ vanishes automatically in two dimensions.
3. Now consider the two-dimensional de Sitter metric

$$
d s^{2}=-d u^{2}+\cosh ^{2} u d \varphi^{2}
$$

where $\infty-<u<\infty$ and $0 \leq \varphi \leq 2 \pi$. Using your results from question 4 of exercise sheet 2 , compute the one independent component of the Riemann tensor, say $R_{u \varphi u \varphi}$.
4. Compute the Ricci scalar $R$ of two-dimensional de Sitter metric and show that indeed $R_{a b c d}=g_{a c} g_{b d}-g_{a d} g_{b c}$.

## 3. Newtonian Limit of Geodesic Deviation

The relative acceleration of a one-parameter family of geodesics $x_{s}^{a}(\tau)=x^{a}(\tau, s)$ is determined by the geodesic deviation equation

$$
\frac{D^{2} X^{d}}{D \tau^{2}}=-R_{a b c}^{d} V^{a} X^{b} V^{c}
$$

where

$$
V^{a}=\frac{d x^{a}}{d \tau} \quad X^{b}=\frac{d x^{b}}{d s}
$$

and $D / D \tau=V^{a} \nabla_{a}$ is the covariant derivative in the direction of $V^{a}$.
In the Newtonian limit, show that this reduces to

$$
\frac{\partial^{2} X^{i}}{\partial t^{2}}=-\sum_{j} \frac{\partial^{2} \phi}{\partial x^{i} \partial x^{j}} X^{j}
$$

where $\phi$ is the gravitational potential.

## 4. Schwarzschild Solution

Everyone should check that the Schwarzschild solution satisfies the vacuum field equations $R_{a b}=0$. You should do this carefully, in your own time, and it is not neccessary to hand this calculation in.

## 5. Circular Orbits

The lagrangian for affinely parametrized geodesics in the Schwarzschild solution is

$$
L=-(1-2 M / r) \dot{t}^{2}+(1-2 M / r)^{-1} \dot{r}^{2}+r^{2}\left(\dot{\theta}^{2}+\sin ^{2} \theta \dot{\phi}^{2}\right)
$$

Show that

$$
E=(1-2 M / r) \dot{t} \quad J=r^{2} \sin ^{2} \theta \dot{\phi}
$$

are conserved. What symmetries do they correspond to? Explain why the lagrangian itself is conserved and why $L=-1$ for time-like geodesics parametrized by the proper time. Show that without loss of generality, you can always restrict attention to time-like geodesics lying in the equatorial plane, $\theta=\pi / 2$. Now consider a circular orbit in the equatorial plane at coordinate distance $R$. Show that this is a time-like geodesic if $R>3 M$ and find the conserved quantities $(E, J)$.

