

Problem Sheet 3, General Relativity 2, HT 2018

1. Let K^a be the null vector with components $(1, 0, 0, 1)$, and let e_{ab} such that

$$\left(e_{ab} - \frac{1}{2} \eta_{ab} e \right) K^b = 0$$

where $e = \eta^{ab} e_{ab}$ (as in the lectures for chapter 2). Show that a 4-vector λ^a can be found so that $\tilde{e}_{ab} = e_{ab} + K_a \lambda_b + \lambda_a K_b$ is given by

$$\tilde{e}_{ab} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & A & B & 0 \\ 0 & B & -A & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}.$$

[Hint: first show that λ^a can be chosen to make $\tilde{e}_{0a} = 0$, then deduce that \tilde{e} must vanish, etc....]

2. A massive particle crosses $r = 2M$ in the Schwarzschild metric following some incoming radial time-like path. Show that it arrives at $r = 0$ after an elapsed proper time Δs which can be no greater than πM .

[As a first step, what class of paths maximise proper time?]

3. Reproduce the figure in Section 3.3 (lecture 12) for the Kruskal-Szekeres spacetime. Show a radial space-like geodesic with $E = 0$, and a radial time-like geodesic with $E < 1$. What does this condition tell you about \dot{r} , and therefore r ? Where on this figure are there radial time-like geodesics with $E = 0$?

4. **How to recognize a null hypersurface:** A null hypersurface is a surface Σ with normal n^μ which is null. In some space-time, the scalar function S has the property that $S = 0$ is a null hypersurface. S is used as a coordinate, $S = x^0$, in a coordinate system (x^0, x^1, x^2, x^3) . Show that the component g^{00} of the contravariant metric vanishes at $S = 0$. Decompose the covariant metric into blocks as

$$g_{ab} = \begin{pmatrix} V & \mathbf{Y}^t \\ \mathbf{Y} & A \end{pmatrix}$$

where A is a (symmetric) 3×3 matrix, \mathbf{Y} a 3 component column vector and V a function.

Show that $\det A = 0$ at $S = 0$.

[So this is one (quick) way to recognize a null hypersurface. For the proof consider the identity $g^{ac} g_{cb} = \delta_b^a$.]

For the Schwarzschild metric, use this to show that the surface $S = r - 2M = 0$ is null in the Eddington- Finkelstein metrics.

5. **Static versus Stationary:** A space-time with a time-like Killing vector K^a is said to be *static* if K^a is hypersurface-orthogonal (recall that this means $K_{[a}\nabla_b K_{c]} = 0$) and stationary otherwise. The Kerr metric in the Boyer-Lindquist coordinates is

$$ds^2 = - \left(1 - \frac{2Mr}{\Sigma}\right) dt^2 - \frac{4Mar}{\Sigma} \sin^2 \theta d\phi dt \\ + \frac{1}{\Sigma} \sin^2 \theta (\Delta \Sigma + 2Mr(r^2 + a^2)) d\phi^2 + \Sigma \left(\frac{1}{\Delta} dr^2 + d\theta^2\right)$$

where $\Sigma = r^2 + a^2 \cos^2 \theta$ and $\Delta = r^2 - 2Mr + a^2$. Show that this metric is not static unless $J = Ma = 0$ (when of course it reduces to the Schwarzschild metric).

6. Show that a vector field \mathbf{K} in flat 3-space is hypersurface orthogonal (HSO) if and only if

$$\mathbf{K} \cdot \nabla \wedge \mathbf{K} = 0 .$$

If $\mathbf{K} = (y, -x, f(r))$ where $r^2 = x^2 + y^2$, for what f is \mathbf{K} HSO? What do the integral curves (equivalently ‘streamlines’) of \mathbf{K} look like?

If $\mathbf{K} = (\cos g(z), \sin g(z), 0)$ where $g' \neq 0$, show that \mathbf{K} is never HSO and that the integral curves of \mathbf{K} are straight lines.

7. The Kerr metric in Kerr-Eddington coordinates is

$$ds^2 = -dT^2 + dr^2 - 2a \sin^2 \theta dr d\Phi + \Sigma d\theta^2 + (r^2 + a^2) \sin^2 \theta d\Phi^2 \\ + \frac{2Mr}{\Sigma} (dT - a \sin^2 \theta d\Phi + dr)^2$$

where T and Φ are defined by the equations

$$dT = dt + \frac{2Mr}{\Delta} dr , \quad d\Phi = d\phi + \frac{a}{\Delta} dr ,$$

and

$$\Delta = r^2 - 2Mr + a^2 , \quad \Sigma = r^2 + a^2 \cos^2 \theta .$$

Use the method in question 4 with this metric to show that $r = r_+$ and $r = r_-$ are null hypersurfaces in the Kerr metric (recall that r_{\pm} are the values of r for which $\Delta = 0$).